## Disproof of Bell's Theorem Illuminating the Illusion of Entanglement

Second Edition

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BrownWalker Press Boca Raton

## Appendix 9: Simulation of EPR-Bohm Correlation



Figure 9.1: The measurement results  $\mathscr{A}(\mathbf{a}; \mathbf{e}_o, \theta_o)$  and  $\mathscr{B}(\mathbf{b}; \mathbf{e}_o, \theta_o)$  are deterministically brought about by the common cause  $(\mathbf{e}_o, \theta_o)$ .

The set of *complete* states  $(\mathbf{e}_o, \theta_o)$  of the system is defined by

$$\Lambda := \left\{ (\mathbf{e}_o, \, \theta_o) \, \middle| \, |\cos(\eta_{\mathbf{x}\mathbf{e}_o})| \ge \frac{1}{2} \, \sin^2(\theta_o) \quad \forall \, \mathbf{x} \in T_p S^3 \right\}, \quad (A.9.1)$$

where  $\mathbf{e}_o \in \mathbb{R}^3$  is a random vector,  $\theta_o \in [0, \pi/2]$  is a random scalar, and each vector  $\mathbf{x}$  specifies a *different* 2-sphere within the 3-sphere. Note that there exist no states for which  $|\cos(\eta_{\mathbf{x}\mathbf{e}_o})| < \frac{1}{2} \sin^2(\theta_o)$ for any  $\mathbf{x} \in T_p S^3$ . The two measurement results are then defined by

$$\mathscr{A}(\mathbf{a}; \mathbf{e}_{o}, \theta_{o}) = \begin{cases} sign\{-\cos(\eta_{\mathbf{a}\mathbf{e}_{o}})\} & \text{if } |\cos(\eta_{\mathbf{a}\mathbf{e}_{o}})| \ge \frac{1}{2}\sin^{2}(\theta_{o}) \\ 0 & \text{if } |\cos(\eta_{\mathbf{a}\mathbf{e}_{o}})| < \frac{1}{2}\sin^{2}(\theta_{o}) \end{cases}$$
(A.9.2)

and

$$\mathscr{B}(\mathbf{b}; \mathbf{e}_{o}, \theta_{o}) = \begin{cases} sign\{+\cos(\eta_{\mathbf{b}\mathbf{e}_{o}})\} & \text{if } |\cos(\eta_{\mathbf{b}\mathbf{e}_{o}})| \geq \frac{1}{2}\sin^{2}(\theta_{o}) \\ 0 & \text{if } |\cos(\eta_{\mathbf{b}\mathbf{e}_{o}})| < \frac{1}{2}\sin^{2}(\theta_{o}), \end{cases}$$
(A.9.3)

with *freely chosen* **a** and **b**. Evidently, the averages of the results  $\mathscr{A}$  and  $\mathscr{B}$  vanish, but the probabilities of their joint detections do not:

$$P^{++} = P\{\mathscr{A} = +1, \ \mathscr{B} = +1 \mid \eta_{\mathbf{ab}}\} = \frac{1}{2}\sin^2\left(\frac{\eta_{\mathbf{ab}}}{2}\right), \quad (A.9.4)$$

$$P^{--} = P\{\mathscr{A} = -1, \ \mathscr{B} = -1 \mid \eta_{\mathbf{ab}}\} = \frac{1}{2}\sin^2\left(\frac{\eta_{\mathbf{ab}}}{2}\right), \qquad (A.9.5)$$

$$P^{-+} = P\{\mathscr{A} = -1, \ \mathscr{B} = +1 \mid \eta_{\mathbf{ab}}\} = \frac{1}{2}\cos^2\left(\frac{\eta_{\mathbf{ab}}}{2}\right), \quad (A.9.6)$$

$$P^{+-} = P\{\mathscr{A} = +1, \ \mathscr{B} = -1 \mid \eta_{\mathbf{ab}}\} = \frac{1}{2}\cos^2\left(\frac{\eta_{\mathbf{ab}}}{2}\right).$$
 (A.9.7)



Figure 9.2: An event-by-event simulation of the correlation between simultaneously occurring measurement events  $\mathscr{A} = \pm 1$  and  $\mathscr{B} = \pm 1$ within a parallelized 3-sphere. The code for this simulation is written by Michel Fodje, in Python. Along with other relevant information, it can be downloaded from https://github.com/minkwe/epr-simple/.

From these joint probabilities (which are exactly those predicted by quantum mechanics) the correlation can now be computed as follows:

$$\begin{aligned} \mathcal{E}(\mathbf{a}, \mathbf{b}) &= \lim_{n \gg 1} \left[ \frac{1}{n} \sum_{i=1}^{n} \mathscr{A}(\mathbf{a}; \mathbf{e}_{o}^{i}, \theta_{o}^{i}) \mathscr{B}(\mathbf{b}; \mathbf{e}_{o}^{i}, \theta_{o}^{i}) \right] \\ &= \frac{P^{++} + P^{--} - P^{-+} - P^{+-}}{P^{++} + P^{--} + P^{-+} + P^{+-}} \\ &= \frac{1}{2} \sin^{2} \left( \frac{\eta_{\mathbf{a}\mathbf{b}}}{2} \right) + \frac{1}{2} \sin^{2} \left( \frac{\eta_{\mathbf{a}\mathbf{b}}}{2} \right) \\ &\quad - \frac{1}{2} \cos^{2} \left( \frac{\eta_{\mathbf{a}\mathbf{b}}}{2} \right) - \frac{1}{2} \cos^{2} \left( \frac{\eta_{\mathbf{a}\mathbf{b}}}{2} \right) \\ &= -\cos \eta_{\mathbf{a}\mathbf{b}} \,. \end{aligned}$$
(A.9.8)

These are then the strong correlations predicted by our local model<sup>1</sup>.

<sup>&</sup>lt;sup>1</sup> The original simulation by Michel Fodje confirming the above results has been translated by John Reed from Python to Mathematica. It can be found in PDF format at this page: http://libertesphilosophica.info/Minkwe\_Sim\_J\_Reed.pdf.